# A theoretical approach to Shor's Algorithm and Quantum Bits

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# 1- Introduction to Qubits

## Introduction



$$\langle \cdot, \cdot \rangle : V^2 \longrightarrow \mathbb{C},$$

that is

sesquilinear:

$$\begin{split} &\langle \lambda x + \mu x', y \rangle = \bar{\lambda} \langle x, y \rangle + \bar{\mu} \langle x', y \rangle \\ &\langle x, \lambda y + \mu y' \rangle = \lambda \langle x, y \rangle + \mu \langle x, y' \rangle \end{split}$$

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symmetric: ⟨x, y⟩ = ⟨y, x⟩
positive: ⟨x, x⟩ > 0 if x ≠ 0.
We define also for x ∈ V the norm ||x|| = √⟨x, x⟩.
Let V, W<sup>1</sup> be hermitian spaces, f : V → W is a unitary morphism if it is C-linear and

$$\langle f(x), f(y) \rangle = \langle x, y \rangle \qquad \forall x, y \in V.$$

<sup>1</sup>actually we will consider always V = W in this presentation.



#### Definition

An *n*-qubit is a vector of norm 1 in a hermitian space  $V \cong \mathbb{C}^{2^n} = (\mathbb{C}^2)^{\otimes n}$ . The set of *n*-qubits is denoted by  $Q_n$ .

In the case 
$$n = 1, 2$$
:  
**1**-qubits:  $|0\rangle := \begin{pmatrix} 1\\0 \end{pmatrix}$ ,  $|1\rangle := \begin{pmatrix} 0\\1 \end{pmatrix}$ ;  
**2**-qubits:  $|00\rangle := \begin{pmatrix} 1\\0\\0\\0 \end{pmatrix}$ ,  $|01\rangle := \begin{pmatrix} 0\\1\\0\\0 \end{pmatrix}$ ,  $|10\rangle := \begin{pmatrix} 0\\0\\1\\0 \end{pmatrix}$ ,  $|11\rangle := \begin{pmatrix} 0\\0\\0\\1 \end{pmatrix}$ 



■ Let  $x, n \in \mathbb{N}$  with  $2^n > x$ , then the  $x^{\text{th}}$  vector of the n-qubits basis is represented as

$$|x\rangle_n = |x_{n-1}x_{n-2}\dots x_0\rangle = |x_{n-1}\rangle \otimes |x_{n-2}\rangle \dots \otimes |x_0\rangle,$$

where  $x = \sum_{j=0}^{n-1} 2^{j} x_{j}$ .

A generic *n*-qubit is represented as a *superposition* 

$$|\psi\rangle_n = \sum_{x=0}^{2^n-1} \alpha_x \, |x\rangle_n \text{ with } \sum_{x=0}^{2^n-1} |\alpha_x|^2 = 1.$$

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Let n, m be positive integers, there's a bilinear map

$$Q_n \times Q_m \longrightarrow Q_{n+m}$$
$$(|\psi\rangle_n, |\phi\rangle_m) \longmapsto |\psi\rangle_n \otimes |\phi\rangle_m.$$

As an example

$$\begin{pmatrix} 0\\0\\1\\0 \end{pmatrix} = |10\rangle = |1\rangle \otimes |0\rangle = \begin{pmatrix} 0\\1 \end{pmatrix} \otimes \begin{pmatrix} 1\\0 \end{pmatrix} = \begin{pmatrix} 0\begin{pmatrix}1\\0\\1\\1 \end{pmatrix} \\ 1\begin{pmatrix}1\\0 \end{pmatrix} \end{pmatrix}$$



Notice that a 2-qubit is not always given by two 1-qubits. As an example

$$\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \neq |\psi\rangle_1 \otimes |\phi\rangle_1 \qquad \forall |\psi\rangle_1, |\phi\rangle_1 \in Q_1.$$

However, we have clearly that  $Q_2$  is spanned by 2-qubits that are given by two 1-qubits.

# Quantum gates

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Gates acting on  $Q_1$ 

Hadamard NOT Phase Shift

$$\begin{split} \mathbf{H} : & |x\rangle \mapsto \frac{1}{\sqrt{2}} (|0\rangle + (-1)^{x} |1\rangle \\ \mathbf{X} : & |x\rangle \mapsto |x \oplus 1\rangle \\ \mathbf{R}_{n} : & |x\rangle \mapsto e^{\frac{2\pi i x}{2^{n}}} |x\rangle \end{split}$$

Matrix representations

$$\mathbf{H}: \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1\\ 1 & -1 \end{pmatrix} \qquad \mathbf{X}: \begin{pmatrix} 0 & 1\\ 1 & 0 \end{pmatrix} \qquad \mathbf{R}_n: \begin{pmatrix} 1 & 0\\ 0 & e^{\frac{2\pi i x}{2^n}} \end{pmatrix}$$

Circuit notation



# Quantum gates



Gates acting on  $Q_2$ 

Controlled NOT SWAP  $\begin{array}{l} \mathbf{C} \, \mathbf{X} : \, |xy\rangle \mapsto |x,y \oplus x\rangle \\ \mathbf{SWAP} : \, |xy\rangle \mapsto |yx\rangle \end{array}$ 

Matrix representations

$$\mathbf{CX} : \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix} \qquad \mathbf{SWAP} : \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

Circuit notation





Let  $f:\{0,1\}^n \rightarrow \{0,1\}^m,$  up to special cases we cannot define a gate acting as

$$\begin{aligned} \mathbf{U}_f : \ Q_n \longrightarrow Q_m \\ |x\rangle_n \longmapsto |f(x)\rangle_n \end{aligned}$$

because is not generally a unitary transformation from a space to itself. Then we consider

$$\begin{aligned} \mathbf{U}_{f} : Q_{n+m} &\longrightarrow Q_{n+m} \\ |x\rangle_{n} \otimes |y\rangle_{m} &\mapsto |x\rangle_{n} \otimes |y \oplus f(x)\rangle_{m} . \end{aligned}$$

# Measuring Qubits

Let  $|\psi\rangle_n$  be a superposition of *n*-qubits. If we want to measure the k first qubits, we write

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$$|\psi\rangle_n = \sum_{x=0}^{2^k-1} |x\rangle_k \otimes |\psi_x\rangle_{n-k} \,.$$

The outcome of the measure is x and the quantum state is left to

$$\frac{|x\rangle_k \otimes |\psi_x\rangle_{n-k}}{||\psi_x\rangle_{n-k}||},$$

with probability

 $|||\psi_x\rangle_{n-k}||^2.$ 

Circuit notation



# 2- Quantum Fourier Transform over



Quantum Fourier Transform over  $\mathbb{Z}_{2^n\mathbb{Z}}$   $\bigcirc$  **EETelsy** 

#### Definition

Let x be a an integer in  $\{0, \ldots, 2^n - 1\}$ , we define the *Quantum* Fourier Transform over  $\mathbb{Z}_{2^n\mathbb{Z}}$  of the *n*-qubit  $|x\rangle_n$  as

$$\mathsf{QFT}_n(|x\rangle_n) = \frac{1}{\sqrt{2^n}} \sum_{y=0}^{2^n-1} e^{\frac{2\pi xy}{2^n}} |y\rangle_n.$$

Let  $w_k$  be  $e^{\frac{2\pi i}{2^k}}$ , then we will need later also the equality:

$$\mathsf{QFT}_n(|x\rangle_n) = \frac{|0\rangle + w_1^x |1\rangle}{\sqrt{2}} \otimes \frac{|0\rangle + w_2^x |1\rangle}{\sqrt{2}} \otimes \cdots \otimes \frac{|0\rangle + w_n^x |1\rangle}{\sqrt{2}}.$$
(1)

The quantum circuit which performs the Quantum Fourier Transform is constructed out of Hadamard and Controlled Phase Shift gates.

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# 3- Shor's Algorithm

Let *N* be a positive integer and *a* be an integer such that gcd(a, N) = 1, then Shor's algorithm aim is to find the period of the function

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 $f(x) = a^x \mod N,$ 

with time complexity polynomial in  $\log_2 N$ .

The algorithm is divided in a main quantum part and a classical post processing. The interpretation of the quantum part is the subject of this presentation.

# 3.1- Hidden Subgroup Problem



#### Problem

Let G be a finitely generated group and X be a set. Given a function  $f : G \to X$  such that there exists a subgroup H < G with the following property

$$f(g) = f(g') \Leftrightarrow g' = gh \ \exists h \in H,$$

find a generating set for H.

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Let G be a group, a *character of* G is a group homomorphism  $\chi: G \to \mathbb{C}^*$ . The set  $\hat{G}$  of characters of G is called the *dual group of* G.

Indeed, the set  $\hat{G}$ , equipped with

$$\hat{G} \times \hat{G} \longrightarrow \hat{G}$$
  
 $(\chi_1, \chi_2) \longmapsto \chi_1 \chi_2 : g \mapsto \chi_1(g) \chi_2(g),$ 

is a group.

From now on, the group G will be a **finite abelian** group. In this particular case we have  $\hat{G} \cong G$ , however the isomorphism is not canonical.

Let  $f : G \to X$ , in this general context the *Quantum Fourier Transform* considered is a gate acting in the following way.

$$\mathsf{QFT}\left(\frac{1}{\sqrt{|G|}}\sum_{g\in G}|g\rangle\otimes\left|f(g)\right\rangle\right)=\frac{1}{\sqrt{|G|}}\sum_{\chi\in\hat{G}}|\chi\rangle\otimes\left|\hat{f}(\chi)\right\rangle,$$

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where

$$\left|\hat{f}(\chi)\right\rangle = \frac{1}{\sqrt{|G|}} \sum_{g \in G} \chi(g) \left|f(g)\right\rangle.$$

Of course, it can be proved that if  $G = \mathbb{Z}_{2^n \mathbb{Z}}$  then applying **QFT**<sub>n</sub> to the register  $|g\rangle$  gives the same result.



Given a function f with the assumptions of HSP, this quantum circuit returns a uniformly distributed  $\chi \in \widehat{G_{H}}$ , where  $\widehat{G_{H}}$  is viewed as the subset of  $\hat{G}$  acting trivial on H.





**1** The gate **US** sends  $|0\rangle^{\otimes n}$  to the uniform superposition

$$|0\rangle^{\otimes n} \longmapsto \frac{1}{\sqrt{|G|}} \sum_{x \in G} |g\rangle.$$

**2** The gate  $U_f$  acts as defined before.

$$rac{1}{\sqrt{|G|}} \sum_{x \in G} \ket{g} \otimes \ket{0}^{\otimes m} \longmapsto rac{1}{\sqrt{|G|}} \sum_{x \in G} \ket{g} \otimes ig| f(g) ig
angle$$



#### 3 QFT gives

$$\begin{split} \frac{1}{\sqrt{|G|}} \sum_{x \in G} |g\rangle \otimes |f(g)\rangle &\mapsto \frac{1}{|G|} \sum_{\chi \in \hat{G}} |\chi\rangle \otimes \Big(\sum_{g \in G} \chi(g) |f(g)\rangle\Big) \\ &= \frac{1}{|G/H|} \sum_{\substack{\chi \in \hat{G} \\ \chi|_{H} = 1}} |\chi\rangle \otimes \sum_{g \in G/H} \chi(g) |f(g)\rangle \end{split}$$

**4** The outcome of the measure is  $\chi \in \widehat{G_{H}}$  with probability

$$\left\|\frac{1}{\left|\overline{G}_{/H}\right|}\sum_{g\in\overline{G}_{/H}}\chi(g)\left|f(g)\right\rangle\right\|^{2}=\frac{1}{\left|\overline{G}_{/H}\right|}$$

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Implementing Shor's algorithm, to find the order r of a modulo N, requires the following setting:



where

$$f: G \to \{0, \dots, N-1\}$$
$$x \longmapsto a^x \mod N,$$

with 
$$G = \mathbb{Z}_{2^n \mathbb{Z}}$$
 and  $H = \langle r \rangle \leq G$ .



Observe that, from a theoretical point of view, the previous setting is well defined if the period r divides  $2^n$ . Clearly, this is not always the case and a classical post processing is generally needed to recover r with good probability.

This is the main reason why Shor's algorithm is a probabilistic algorithm.

Following the same line as factoring, Shor provides a solution to discrete logarithm problem (DLP) in a cyclic group  $C = \langle g \rangle$  of order M. Let  $x \in C$ , the HSP setting to find  $y \in \mathbb{Z}/_{M\mathbb{Z}}$  such that  $g^y = x$  is described below.

The group is

$$G = \mathbb{Z}_{M\mathbb{Z}} \times \mathbb{Z}_{M\mathbb{Z}}$$

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The function is

$$f: G \longrightarrow C$$
$$(a,b) \longmapsto g^{a} x^{-b}$$

The hidden subgroup is

$$H = \langle y, 1 \rangle \leq G$$

# 3.2- Quantum Phase Estimation



#### Problem

Let **U** be a unitary transformation. Given an eigenstate  $|\psi\rangle$  of **U** find the phase  $\theta \in [0, 1)$  describing its eigenvalue

$$\mathbf{U} \left| \psi \right\rangle = e^{2\pi i \theta} \left| \psi \right\rangle.$$



We point out the following main ingredient.

$$\mathbf{C} \mathbf{U} \frac{|0\rangle + |1\rangle}{\sqrt{2}} \otimes |\psi\rangle = \frac{1}{\sqrt{2}} \Big( |0\rangle |\psi\rangle + e^{2\pi i \theta} |1\rangle |\psi\rangle \Big)$$
$$= \frac{|0\rangle + e^{2\pi i \theta} |1\rangle}{\sqrt{2}} \otimes |\psi\rangle$$

In this notation  ${\bf C}\,{\bf U}$  can be interpreted as a gate acting just on the first qubit since the last part  $|\psi\rangle$  is fixed.

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Given a unitary transformation U acting on m-qubits and an its eigenstate  $|\psi\rangle_m$ , this quantum circuit computes  $2^n\theta$  where  $\theta$  is the phase of the corresponding eigenvalue.





**1** The gate in the middle sends  $|x\rangle_n |\psi\rangle_m$  to  $|x\rangle_n \mathbf{U}^x |\psi\rangle_m$ . It is constructed out of  $2^j$  gates **U** acting on the register  $|\psi\rangle$  controlled by the *j*-th qubit for all *j*'s.



**1** The gate in the middle sends  $|x\rangle_n |\psi\rangle_m$  to  $|x\rangle_n \mathbf{U}^{\times} |\psi\rangle_m$ . It is constructed out of  $2^j$  gates **U** acting on the register  $|\psi\rangle$ controlled by the *j*-th qubit for all *j*'s. As an example, if x = n = 21  $|\psi\rangle_m \equiv U^2$  $|\psi\rangle_{m}$ Hence, previous remark implies that the state in 1 is  $\frac{1}{\sqrt{2^n}}(|0\rangle+|1\rangle)^{\otimes n}\otimes|\psi\rangle_m\mapsto\frac{|0\rangle+w_1^{2^n\theta}|1\rangle}{\sqrt{2}}\otimes\cdots\frac{|0\rangle+w_n^{2^n\theta}|1\rangle}{\sqrt{2}}\otimes|\psi\rangle_m.$ The first register is exactly the Quantum Fourier Transform applied to  $|x\rangle_n = |2^n\theta\rangle_n$ , see (1).

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2 Applying the Inverse Quantum Fourier Transform over  $\mathbb{Z}_{2^n\mathbb{Z}}$  to the first register gives

$$\frac{|0\rangle + w_1^{2^n\theta} |1\rangle}{\sqrt{2}} \otimes \cdots \frac{|0\rangle + w_n^{2^n\theta} |1\rangle}{\sqrt{2}} \mapsto |2^n\theta\rangle_n.$$

This works well if  $2^n\theta$  is an integer, which is not always true. In the general case, the circuit returns an estimation of  $\theta$  which allows us to recover it through a continued fraction argument with good probability.

## Shor as QPE

Implementing Shor's algorithm, to find the order r of a modulo N, requires the following setting:

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$$|0\rangle^{\otimes n=2\lceil \log_2 N\rceil} = H^{\otimes n} = QFT_n^{\dagger}$$

$$\frac{1}{\sqrt{r}} \sum_{s=0}^{r-1} |\psi_s\rangle_m = |1\rangle_{m=\lceil \log_2 N\rceil} = U^{\times}$$

where

$$\mathbf{U}: \left| y \right\rangle_m \longmapsto \left| ay \mod N \right\rangle_m$$

and

$$\left|\psi_{s}\right\rangle_{m} = \frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} e^{-\frac{2\pi i s k}{r}} \left|a^{k} \mod N\right\rangle_{m} \text{ s.t. } \mathbf{U}\left|\psi_{s}\right\rangle_{m} = e^{\frac{2\pi i s}{r}} \left|\psi_{s}\right\rangle_{m}.$$

To avoid any inconvenience in producing  $\left|\psi_{s}\right\rangle_{m}$  for some s, we observe

$$\left|1\right\rangle_{m} = \frac{1}{\sqrt{r}} \sum_{s=0}^{r-1} \left|\psi_{s}\right\rangle_{m}.$$

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Therefore, using  $|1\rangle$  which is a uniform superposition of those eigenstates and reasoning by linearity, the final measure gives

$$\frac{2^n s}{r}$$
,

for s a random integer between 0 and r - 1.

# 4- Breaking RSA



Given N = pq, Alice wants to send a message  $b \in \mathbb{Z}/N\mathbb{Z}^*$  to Bob. Bob's public key is  $c \in \mathbb{Z}/(p-1)(q-1)\mathbb{Z}^*$ , then Alice sends him

$$a \equiv b^c \mod pq.$$

Assume Eve can detect the order r of a, gcd(r, c) = 1 implies that r is also the order of b. Moreover, there exists d such that  $cd \equiv 1 \mod r$ .

$$a^d \equiv b^{cd} \equiv b^{1+mr} \equiv b \mod pq.$$

It can be proved that there's a good probability that the detected period r is even. If so, we have

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$$a^r \equiv 1 \mod pq$$
  $a^{\frac{r}{2}} \not\equiv 1 \mod pq$ .

Assume also that

$$a^{\frac{r}{2}} \not\equiv -1 \mod pq,$$

since

$$(a^{rac{r}{2}}-1)(a^{rac{r}{2}}+1)\equiv 0 \mod pq$$

we conclude

$$\{p,q\} = \{\gcd(a^{rac{r}{2}}-1,N), \gcd(a^{rac{r}{2}}+1,N)\}.$$



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